文章编号: 1000-4750(2001)03-131-05

复合材料大变形任意加筋板单元

章向明,王安稳

(海军工程大学基础部,武汉 430033)

摘 要:本文构造了用于复合材料偏心加筋板状结构大变形分析的任意加筋板单元。在 此模型中,肋骨可放置在板单元内任意位置和任意方向,因而在网格划分时,可不必顾 及肋骨的具体位置。在板和肋骨的基本方程中,引用 Von-Karman 大变形理论计及大变形 的影响,按照 Mindlin-Reissoner 一阶剪切变形理论考虑横向剪切变形。任意加筋板单元具 有和常规板单元相同的自由度。数值计算表明本文的有限元模型具有很高的精度。

关键词:复合材料加筋板;大变形;横向剪切变形;有限元 中图分类号:O242,U663 文献标识码:A

1 引言

复合材料加筋板状结构在桥梁、航空航天、造船工程和海洋结构等方面有着广泛的用 途。复合材料层合板具有很高的比强度和比刚度,这一特性能大大减轻结构的重量。加强 肋骨更能增强加筋板的承载能力。但复合材料加筋板结构固有的复杂性使分析计算变得非 常困难。已有很多学者用有限元方法研究了复合材料板壳结构,但对复合材料加筋板的研 究却很少,对复合材料加筋板大变形行为的研究就更少^[1, 2]。Sinha G 等人对任意加筋壳结 构进行了线弹性分析^[3],在其有限元模型中采用 36 自由度的任意三角形单元模拟壳结构, 肋骨则采用与壳单元相同的位移形函数。





图 1 任意偏心加筋板单元 图 2 肋骨在板单元中的位置 本文构造的任意加筋板单元如图 1 所示。肋骨可放在板单元的任意位置和任意方向,

收稿日期: 1999-11-27; 修改日期: 2000-03-17 作者简介:章向明(1963),男,安徽庐江人,副教授,博士,主要从事力学研究

而不必放在板单元的节点线上,也不必平行总体坐标系 x 轴或 y 轴:在单元网格划分时, 只须考虑便于应力输出和结果观察而不必顾及肋骨的位置,这就给网格划分带来了很大的 灵活性。在这一模型中,盖板连同加强肋骨的整体被视为一个单元——加筋板单元,其单 元刚度将由板的贡献和梁的贡献共同组成,并假定肋骨具有与盖板相同的位移场,因而肋 骨的位移和相应的广义应变就可以用板单元的形状函数及其相应的导数表达。非线性平衡 方程由 Newton-Raphson 增量迭代法^[4]求解。

2 板的切线刚度

在所构造的模型中,8节点 Serendipity 单元和9节点 Largerange 或 Heterosis 单元等二 次等参数板单元可用于对板进行离散化。

广义应变—位移关系由 Von-Karman 假设给出

$$\{\boldsymbol{e}\} = \{\boldsymbol{e}_{x} \quad \boldsymbol{e}_{y} \quad \boldsymbol{g}_{xy} \quad \boldsymbol{k}_{x} \quad \boldsymbol{k}_{y} \quad \boldsymbol{k}_{xy} \quad \boldsymbol{g}_{xz} \quad \boldsymbol{k}_{yz}\}^{T} = \{\boldsymbol{e}_{0}\} + \{\boldsymbol{e}_{l}\} = \sum_{i=1}^{N_{x}} \boldsymbol{B}_{i} \delta_{i}$$
(1)

其中, $\{e_0\}$ 和 $\{e_i\}$ 是应变的线性部分和非线性部分

$$\{\boldsymbol{e}_{0}\} = \{\partial u/\partial x \quad \partial v/\partial y \quad \partial v/\partial x + \partial u/\partial y \quad \partial \boldsymbol{f}_{x}/\partial x \quad \partial \boldsymbol{f}_{y}/\partial y \quad \partial \boldsymbol{f}_{y}/\partial x + \partial \boldsymbol{f}_{x}/\partial y$$
$$\boldsymbol{f}_{x} + \partial w/\partial x \quad \boldsymbol{f}_{y} + \partial w/\partial y\}^{T}$$
$$\{\boldsymbol{e}_{i}\} = \frac{1}{2}\{(\partial w/\partial x)^{2} \quad (\partial w/\partial y)^{2} \quad (\partial w/\partial x)(\partial w/\partial y) \quad 0 \quad 0 \quad 0 \quad 0\}^{T}$$
$$\{\boldsymbol{d}_{i}\} = \{u_{i}, v_{i}, w_{i}, \boldsymbol{f}_{xi}, \boldsymbol{f}_{yi}\}^{T}$$

应变矩阵可分解成无限小变形线性部分B₀和大变形的非线性部分B₁

$$\boldsymbol{B}_{i} = \boldsymbol{B}_{0i} + \boldsymbol{B}_{li} \tag{2}$$

板对单元刚度矩阵的贡献可写成

$$K_{p}^{T} = \overline{K}^{p} + K_{s}^{p} \tag{3}$$

其中, $\overline{K}^{p} = \int B^{T} D_{p} B dA = \int B^{T} D_{p} B |J| dx dh$, $K_{s}^{p} = \int G^{T} \begin{bmatrix} N_{x} & N_{xy} \\ N_{xy} & N_{y} \end{bmatrix} G dA$; N_{x}, N_{xy}, N_{y} 为 当前应力场的应力合力。

3 肋骨的切线刚度

在偏心加筋板单元中, 肋骨可以放在板单元内的任何位置和任意方向(图 2),但其具体位置须给定。肋骨在板单元中的位置由其端点的自然坐标定义,梁的中点定义为 $((\mathbf{x}_1 + \mathbf{x}_3)/2, (\mathbf{h}_1 + \mathbf{h}_3)/2)$,肋骨轴线上任意点在板单元中的位置按插值方式得到

$$\mathbf{x} = \sum_{k=1}^{\infty} N_k(\mathbf{g}) \mathbf{x}_K, \quad \mathbf{h} = \sum_{k=1}^{\infty} N_k(\mathbf{g}) \mathbf{h}_K$$
(4)
式中, g 是梁单元的自然坐标。

假定变形前垂直于板中面和肋骨中面的公共法线在变形后仍保持为直线,则肋骨中面 上任意点的广义位移可由板中面相应点的广义位移表示成

$$\begin{cases} u^{s} \\ v^{s} \\ w^{s} \\ \mathbf{f}_{x}^{s} \\ \mathbf{f}_{y}^{s} \end{cases} = \begin{bmatrix} 1 & 0 & 0 & e & 0 \\ 0 & 1 & 0 & 0 & e \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} u \\ v \\ w \\ \mathbf{f}_{x} \\ \mathbf{f}_{y} \end{bmatrix}$$
(5)

 u^s, v^s 等为肋骨中面上点的广义位移, e 是肋骨的偏心量: u, v等是板中面上对应点的广义 位移。

每个节点 5 个自由度的 c^0 连续二次等参数板单元可用于对板进行离散化, 单元盖板内 的广义位移按常规方式进行插值

$$\{\boldsymbol{u}^{e}\} = \sum_{i=1}^{N} [N_{i}^{e}] \{\boldsymbol{d}_{i}\}$$
(6)

其中, $\{\boldsymbol{d}_i\} = \{u_i, v_i, w_i, \boldsymbol{f}_{xi}, \boldsymbol{f}_{yi}\}^T$, $[N_i^e] = N_i[\boldsymbol{I}]$, N_i 是常规的插值函数; $[\boldsymbol{I}] \ge 5 \times 5$ 的 单位矩阵: u_i, v_i等是板单元的位移节点值。由式(5)、(6)知: 肋骨的位移场可由其所处板 单元的节点自由度和板单元的形状函数表示。

肋骨的应变应在局部坐标系中计算且由 Von-Karman 假设给出

$$\{\boldsymbol{e}^{s}\} = \{\boldsymbol{e}^{s}_{x'} \ \boldsymbol{k}^{s}_{x'} \ \boldsymbol{k}^{s}_{x'y'} \ \boldsymbol{g}^{s}_{x'z}\}^{T} = \{\boldsymbol{e}^{s}_{0}\} + \{\boldsymbol{e}^{s}_{l}\} = \sum_{i=1}^{N} (\boldsymbol{B}^{s}_{0i} + \boldsymbol{B}^{s}_{li})\boldsymbol{d}_{i} = \sum_{i=1}^{N} \boldsymbol{B}^{s}_{i}\boldsymbol{d}_{i}$$
(7)

其中

$$\{\boldsymbol{e}_{0}^{s}\} = \{\partial u^{\prime s} / \partial x^{\prime} \quad \partial \boldsymbol{f}_{x^{\prime}}^{s} / \partial x^{\prime} \quad \partial \boldsymbol{f}_{y^{\prime}}^{s} / \partial x^{\prime} + \partial \boldsymbol{f}_{x^{\prime}}^{s} / \partial y^{\prime} \quad \boldsymbol{f}_{x^{\prime}}^{s} + \partial w^{\prime} / \partial x^{\prime}\}$$
$$\{\boldsymbol{e}_{l}^{s}\} = \frac{1}{2}\{(\partial w^{\prime} / \partial x^{\prime})^{2} \quad 0 \quad 0 \quad 0\}^{T}$$

位移在局部坐标系下的导数将由整体坐标系下的导数变换得到

$$\begin{bmatrix} \frac{\partial u'^{s}}{\partial x'} & \frac{\partial v'^{s}}{\partial x'}\\ \frac{\partial u'^{s}}{\partial u'^{s}} & \frac{\partial v'^{s}}{\partial y'} \end{bmatrix} = [q]^{T} \begin{bmatrix} \frac{\partial u^{s}}{\partial x} & \frac{\partial v^{s}}{\partial x}\\ \frac{\partial u^{s}}{\partial y} & \frac{\partial v^{s}}{\partial y} \end{bmatrix} [q], \begin{bmatrix} \frac{\partial f_{x'}^{s}}{\partial x'} & \frac{\partial f_{y'}^{s}}{\partial x'}\\ \frac{\partial f_{x'}^{s}}{\partial y'} & \frac{\partial f_{y'}^{s}}{\partial y'} \end{bmatrix} = [q]^{T} \begin{bmatrix} \frac{\partial f_{x}^{s}}{\partial x} & \frac{\partial f_{y}^{s}}{\partial x}\\ \frac{\partial f_{x}^{s}}{\partial y} & \frac{\partial f_{y}^{s}}{\partial y} \end{bmatrix} [q]$$

$$(8)$$

$$\frac{\partial w}{\partial x'} = \frac{\partial w}{\partial x} \cos a + \frac{\partial w}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial w}{\partial x} \cos a + \frac{\partial w}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial \psi}{\partial x} \cos a + \frac{\partial w}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial \psi}{\partial x} \cos a + \frac{\partial w}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial \psi}{\partial x} \cos a + \frac{\partial \psi}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial \psi}{\partial x} \cos a + \frac{\partial \psi}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial \psi}{\partial x} \cos a + \frac{\partial \psi}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial \psi}{\partial x} \cos a + \frac{\partial \psi}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial \psi}{\partial x} \cos a + \frac{\partial \psi}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial \psi}{\partial x} \cos a + \frac{\partial \psi}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial \psi}{\partial x} \cos a + \frac{\partial \psi}{\partial y} \sin a$$

$$\frac{\partial \psi}{\partial x'} = \frac{\partial \psi}{\partial x} - \frac{\partial f_{x}}{\partial y} - \frac{\partial f_{x}}{\partial y} - \frac{\partial f_{x}}{\partial y} - \frac{\partial f_{x}}{\partial y} - \frac{\partial \psi}{\partial y} - \frac{\partial f_{x}}{\partial y} - \frac{\partial \psi}{\partial y} - \frac{\partial f_{x}}{\partial y} - \frac{\partial f_{x}}{\partial y} - \frac{\partial f_{x}}{\partial y} - \frac{\partial f_{y}}{\partial y} - \frac$$

Г. .

~ ~ ¹

$$\begin{bmatrix} \frac{\partial u^{s}}{\partial x} & \frac{\partial v^{s}}{\partial x} & \frac{\partial w}{\partial x} & \frac{\partial \mathbf{f}_{x}^{s}}{\partial x} & \frac{\partial \mathbf{f}_{y}^{s}}{\partial x} \\ \frac{\partial u^{s}}{\partial y} & \frac{\partial v^{s}}{\partial y} & \frac{\partial w}{\partial y} & \frac{\partial \mathbf{f}_{x}^{s}}{\partial y} & \frac{\partial \mathbf{f}_{y}^{s}}{\partial y} \end{bmatrix} = J^{-1} \begin{bmatrix} \frac{\partial u^{s}}{\partial \mathbf{x}} & \frac{\partial v^{s}}{\partial \mathbf{x}} & \frac{\partial w}{\partial \mathbf{x}} & \frac{\partial \mathbf{f}_{x}^{s}}{\partial \mathbf{x}} & \frac{\partial \mathbf{f}_{y}}{\partial \mathbf{x}} \\ \frac{\partial u^{s}}{\partial \mathbf{h}} & \frac{\partial v^{s}}{\partial \mathbf{h}} & \frac{\partial w}{\partial \mathbf{h}} & \frac{\partial \mathbf{f}_{x}^{s}}{\partial \mathbf{h}} & \frac{\partial \mathbf{f}_{y}^{s}}{\partial \mathbf{h}} \end{bmatrix}$$
(9)

其中, J是雅可比矩阵。

式(7)中非线性应变可写成

$$\{\boldsymbol{e}_{i}^{s}\} = \frac{1}{2} \begin{bmatrix} \frac{\partial w}{\partial x'} & 0 & 0 \end{bmatrix}^{T} \frac{\partial w}{\partial x'} = \frac{1}{2} \boldsymbol{A}^{s} \boldsymbol{R}^{s}$$
(10)

其中

$$\boldsymbol{R}^{s} = \frac{\partial w}{\partial x'} = \sum_{i=1}^{N} \boldsymbol{G}_{i}^{s} \boldsymbol{d}_{i}$$
(11)

非线性应变矩阵可写成

$$\boldsymbol{B}_{li}^{s} = \boldsymbol{A}^{s} \boldsymbol{G}_{l}^{s} \tag{12}$$

肋骨对单元刚度矩阵的贡献可写成

$$\boldsymbol{K}_{s}^{T} = \overline{\boldsymbol{K}}^{s} + \boldsymbol{K}_{s}^{s}$$
(13)

其中

$$\overline{\boldsymbol{K}}^{s} = \int \boldsymbol{B}_{s}^{T} \boldsymbol{D}_{s} \boldsymbol{B}_{s} b dx' = \int \boldsymbol{B}_{s}^{T} \boldsymbol{D}_{s} \boldsymbol{B}_{s} b \big| \boldsymbol{J}_{s} \big| d\boldsymbol{g}$$
$$\boldsymbol{K}_{\boldsymbol{S}}^{s} = \int \boldsymbol{G}_{s}^{T} N_{x'}^{s} \boldsymbol{G}_{s} b dx' = \int \boldsymbol{G}_{s}^{T} N_{x'}^{s} \boldsymbol{G}_{s} b \big| \boldsymbol{J}_{s} \big| d\boldsymbol{g}$$

层合板的本构方程及考虑 y 方向 Poisson 效应的层合梁的本构方程已由文献[5]导出。

4 数值计算

算例 1: 四边简支各向同性正方形加筋板(图 3)受均布载荷 q=1.0psi 作用。材料性质为: $E=17 \times 10^6$, n=0.3。采用两种单元网格进行计算,网格1只是用来说明所构造的单元能够 处理任意方向的肋骨布置情形的能力,两种网格计算得到加筋板中心点的挠度基本相同, 计算结果列于表1中;网格2用于计算沿板中心线的挠度分布,对于同心和偏心加筋板沿 板中心线的挠度分布曲线如图4所示。在图4和表1中,由本文计算结果与文献[3]比较可 以看出二者较吻合。



图 3 四边简支正方形加筋板几何尺寸及有限元网格

算例 2: 四边简支两层加筋层合矩形板如图 5 所示。肋骨和盖板的铺设方案均为 $[0^{0}/90^{0}]$,加筋板的几何尺寸为: *a*=254mm,*b*=508mm,*c_x*=6.35mm,*d_x*=12.7mm; 纤维增 强复合材料的性质为: *E*₁=144.8GPa, *E*₂=9.67GPa, *G*₁₂=*G*₁₃=4.14GPa, *G*₂₃=3.45GPa, *n*₁₂=0.3, *r*=1389.23kg/m³; 整块板用 4×4 有限元网格进行离散化,在均布载荷的作用下 计算得到矩形板中心点的非线性挠度曲线与文献[2]的比较示于图 5 中,可见二者符合得很 好。

134

5 结论

表1

本文构造的复合材料加筋板单元具有与常规板单元相同的自由度,它能适应肋骨的任 意分布,因而在实际操作中容易实施且给网格划分带来很大的灵活性。通过数值计算可知, 该单元用相对稀疏的单元网格和较少的机时就能获得很高的精度。

来源	均布载荷	
	对称加筋	偏心加筋
NASTRAN	0.1490	_
STRUDL		0.4539
文献 3	0.1367	0.4556
本文结果	0.1363	0.4554

正方形加筋板中心点挠度



图 4 四边简支正方形加筋板中心线挠度曲线



图 5 带有 x 方向肋骨的矩形加筋板的几何尺寸及中心点的非线性挠度曲线

参考文献:

- C L Liao and J N Reddy. Analysis of anisotropic stiffened composite laminates using a continuum-based shell element[J]. Comput Struct, 1990, 34(6): 805-815.
- [2] M Kolli and K Chandrashekhara. Non-linear static and dynamic analysis of stiffened laminated plates[J]. Int J Non-linear Mechanics, 1997, 32(1): 89-101.
- [3] G Sinha, A H Sheikh and M Mukhopadyay. A new finite element model for the analysis of arbitrary stiffened shells[J]. Finite Elements in Analysis and Design, 1992, 12: 241-271.
- [4] J A Figueiras and D R J Owen. Analysis of elasto-plastic and geometrically nonlinear anisotopic plates and shells[A]. In: Hinton E, Owen D R J eds. Finite element software for plates and shells[C]. Pineridge press limited, 1984.
 - [5] X M Zhang and A W Wang. Large deformation beam-plate element for eccentrically laminated stiffened plates[A]. In: Xu B Y, Tokuda M and Wang X C eds. Microstructures and Mechanical Properties of New Engineering Materials[C]. International Academic Publisher, Beijing, China, 1999. (下转 130 页)

[5] 匡文起, 等. 结构矩阵分析和程序设计[M]. 高等教育出版社, 1991.

FINITE ELEMENT ANALYSIS OF LARGESCALE SKEW

AQUEDUCT BRIDGE

ZHAO Ping, TANG Ke-dong, LIU Zuo-qiu, CHEN Wen-yi, LI Shu-yao

(Civil Engineering Department, North China Institute of Water Conservancy and Hydropower, Zhengzhou 450045)

Abstract: This paper presents a spatial frame element analysis of large scale skew aqueduct bridge. Comparison is made between skew aqueduct bridges and straight aqueduct bridges. It is found that the internal force distribution of a skew aqueduct bridge is complicated compared with that of a straight aqueduct bridge. A skew aqueduct bridge is characterized by the internal forces of a composite structure.

Key words: aqueduct bridge; skew crossbeam; FEM; internal force

(上接135页)

GEOMETRICALLY NONLINEAR PLATE ELEMENT FOR

ARBITRARILY STIFFENED LAMINATED COMPOSITE PLATES

ZHANG Xiang-ming, WANG An-wen

(Naval University of Engineering, wuhan 430033)

Abstract: A new stiffened plate element is constructed to investigate the geometrically nonlinear behavior of arbitrarily stiffened laminated composite plates. In this model, the stiffener can be placed at anywhere within the plate element and oriented in any direction. The Von-Karman kinematic relations of the plate and stiffener are considered. The transverse shear deformation effect is included by use of Mindlin-Reissner first-order shear deformation theory. This new stiffened plate element has the same degree of freedom as regular plate elements. The results obtained by the proposed method indicate that good agreement with available values can be reached with relatively coarse mesh and less computational time.

Key words: stiffened laminated plate; large deformation; transverse shear; finite element

^[6] C A Brebbia, A J Ferrante. Computational methods for the solution of engineering problems[M]. Third Revised Edition, Pentech Press Limited, 1986.